



Pavol Jozef Šafárik University in Košice

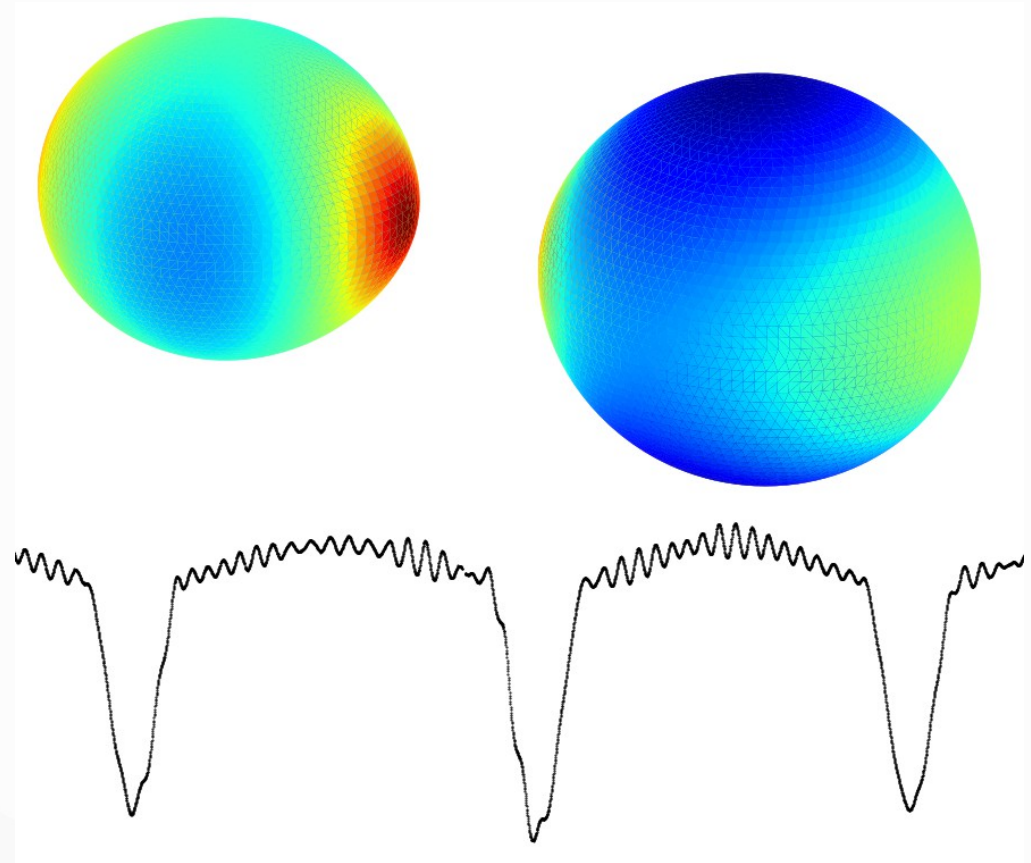
Modelling of eclipsing binaries with pulsating components

Košice 2019

Mgr. Miroslav Fedurco

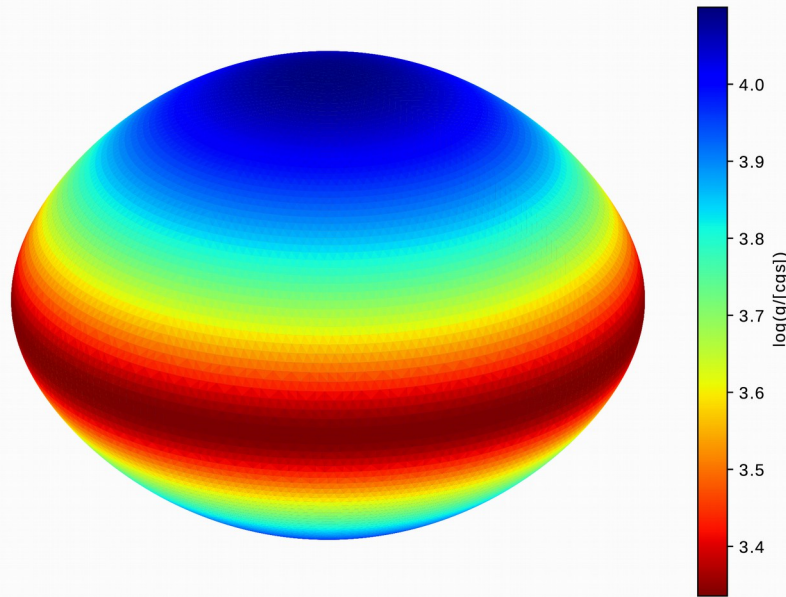
Motivation

- **Large amount of oscillating binaries in data from photometric surveys and space based telescopes.**
- **Eclipsing binaries enable precise derivation of absolute parameters of components.**
- **Great starting point for subsequent asteroseismological analysis.**

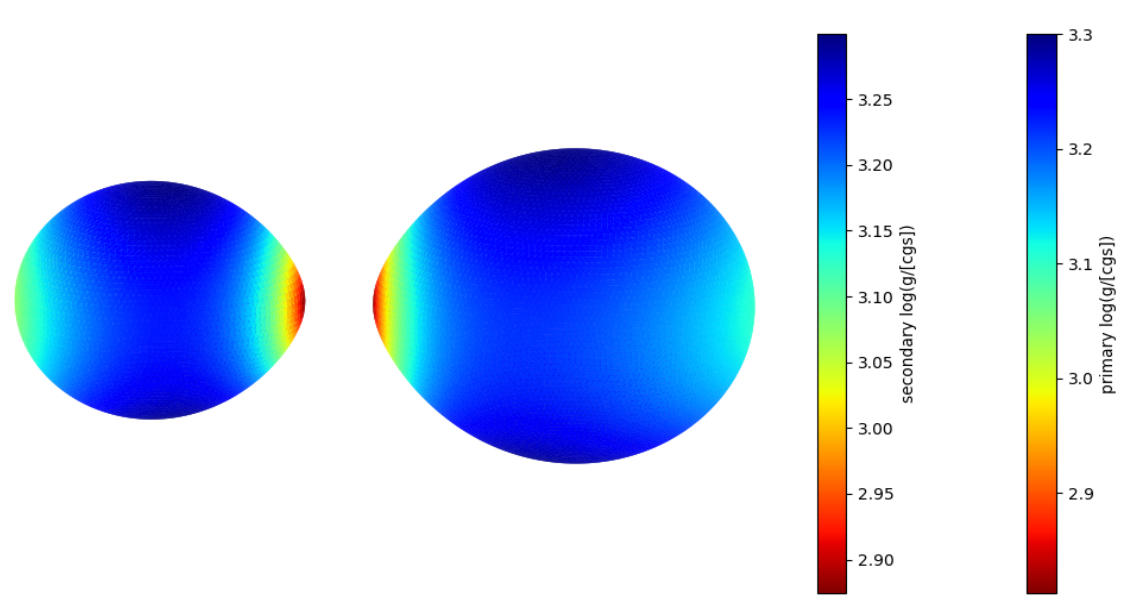


Main objects of interest

- **Rapidly rotating stars**



- **Tidally deformed stars**



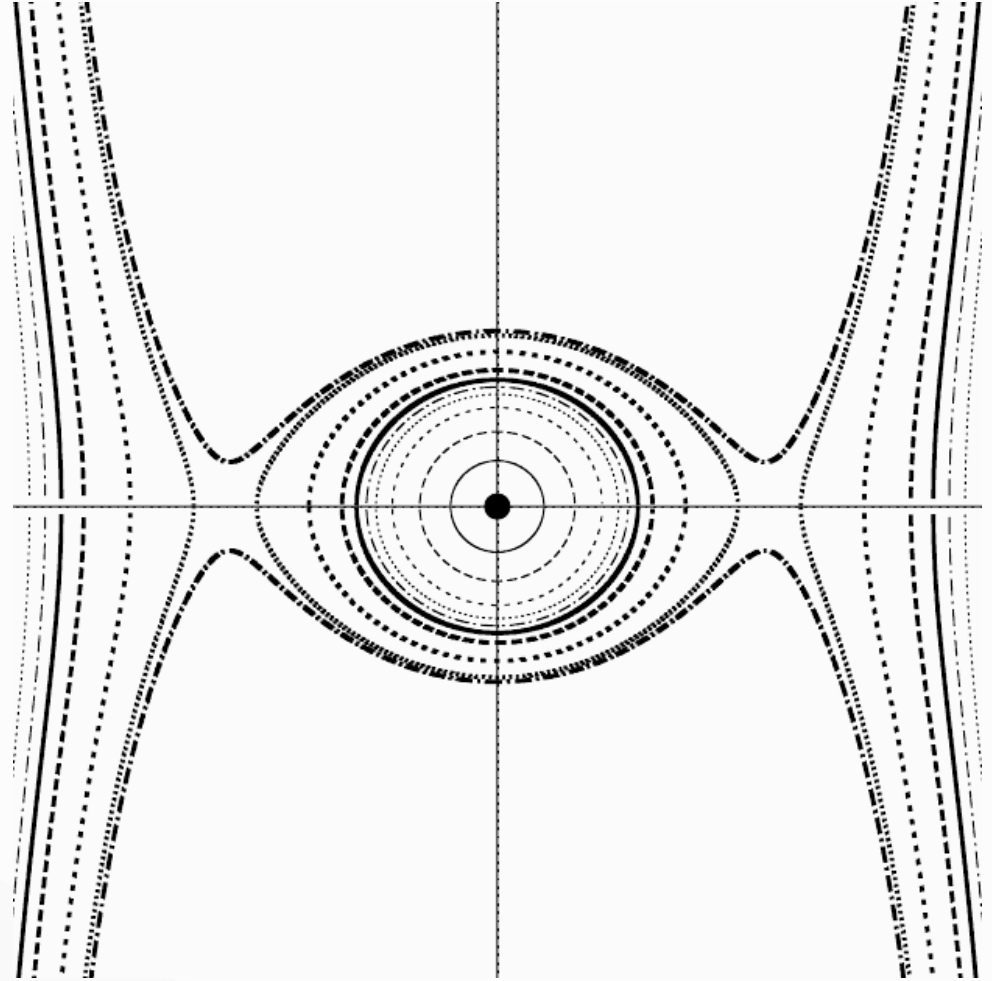
Rapidly rotating stars

- **Combination of gravitational and centrifugal force.**

$$\Psi = G \frac{M}{r} + \frac{1}{2} (r\omega \sin \theta)^2$$

- **Rapid rotation causes equatorial flattening.**
- **Existence of a critical break-up equatorial rotation speed.**

$$v_c = \sqrt[3]{GM\omega}$$

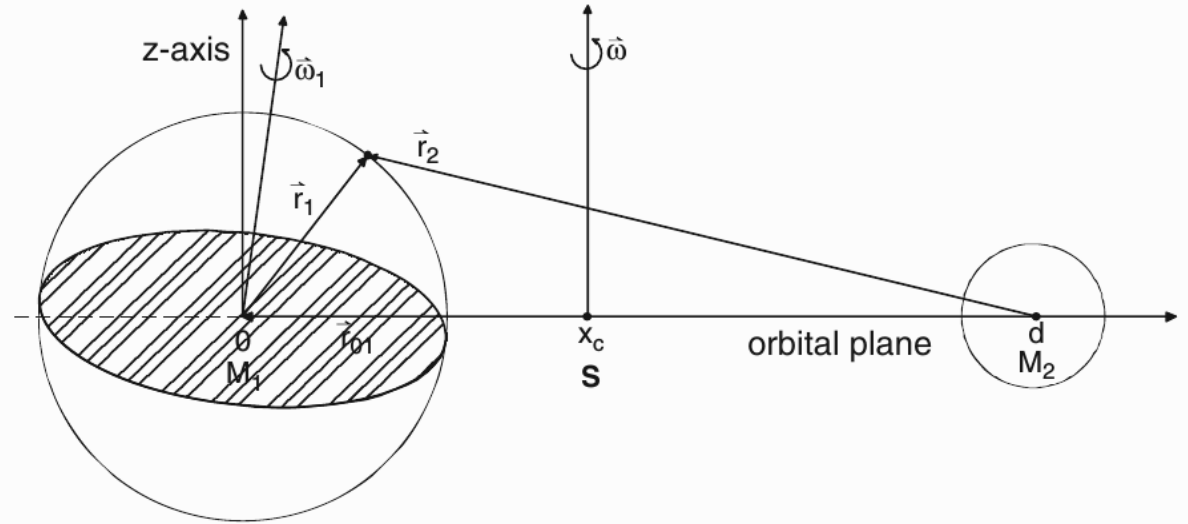


Source: Neslušan (2009)

Roche geometry

- **Restricted 3 body problem**

- Stars regarded as a point masses.
- Uniform rotation.
- Surface in hydrostatic equilibrium.

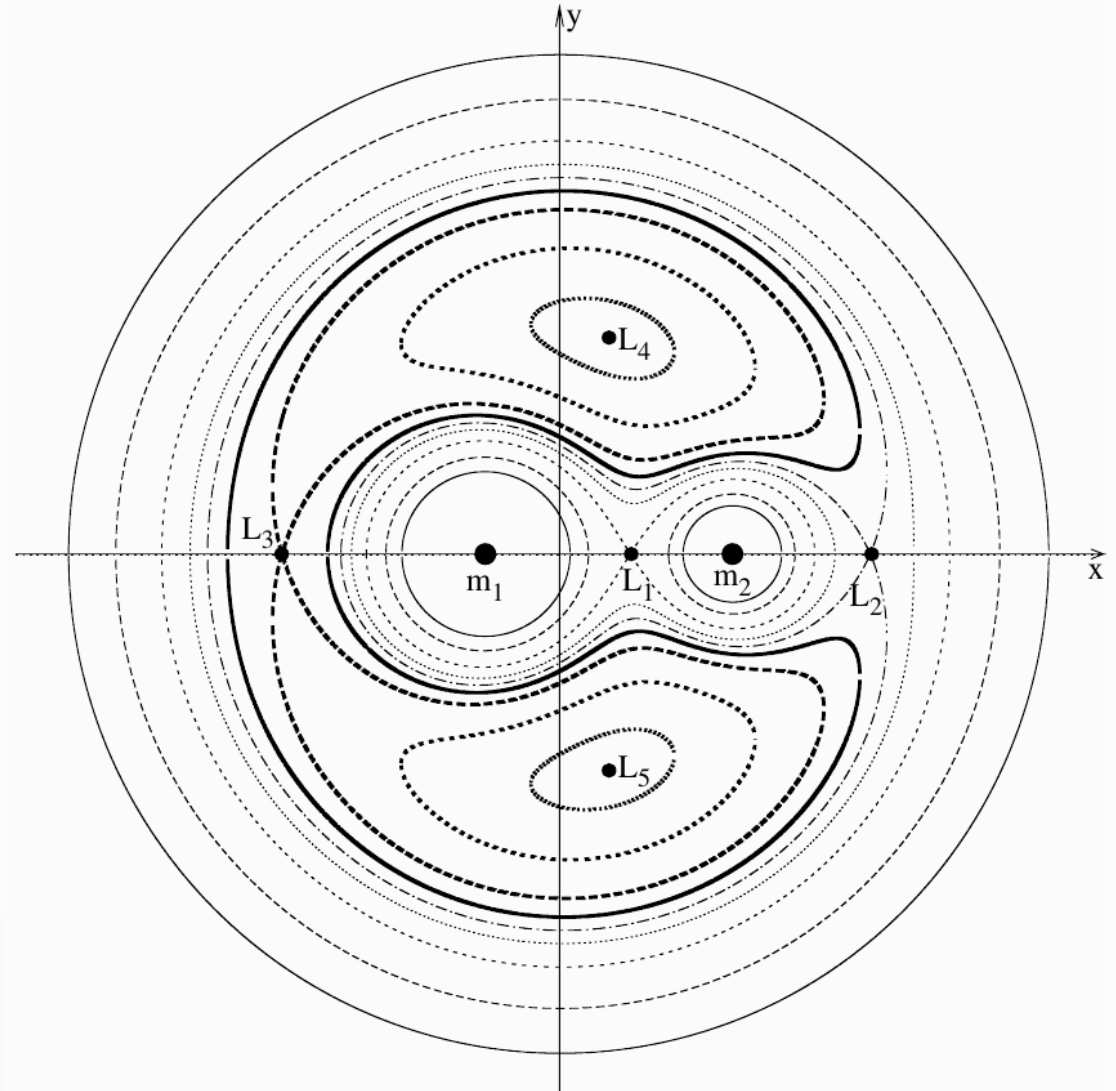


Source: Kallrath & Milone (2009)

$$\Omega(\varrho; q, d) = \frac{1}{\varrho} + q \left[\frac{1}{\sqrt{d^2 - 2d\varrho \cos \varphi \cos \theta + \varrho^2}} + \frac{\varrho \cos \varphi \cos \theta}{d^2} \right] + \frac{1}{2} \left(\frac{\omega_1}{\omega} \right)^2 (q + 1) \varrho^2 (1 - \cos^2 \theta)$$

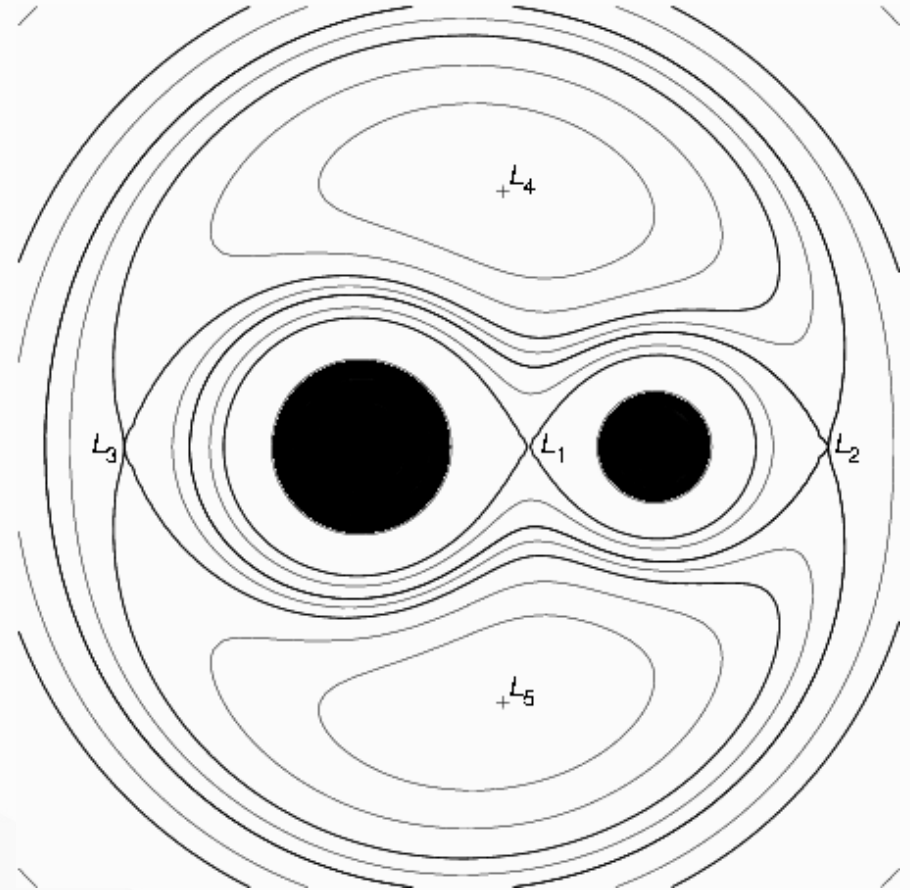
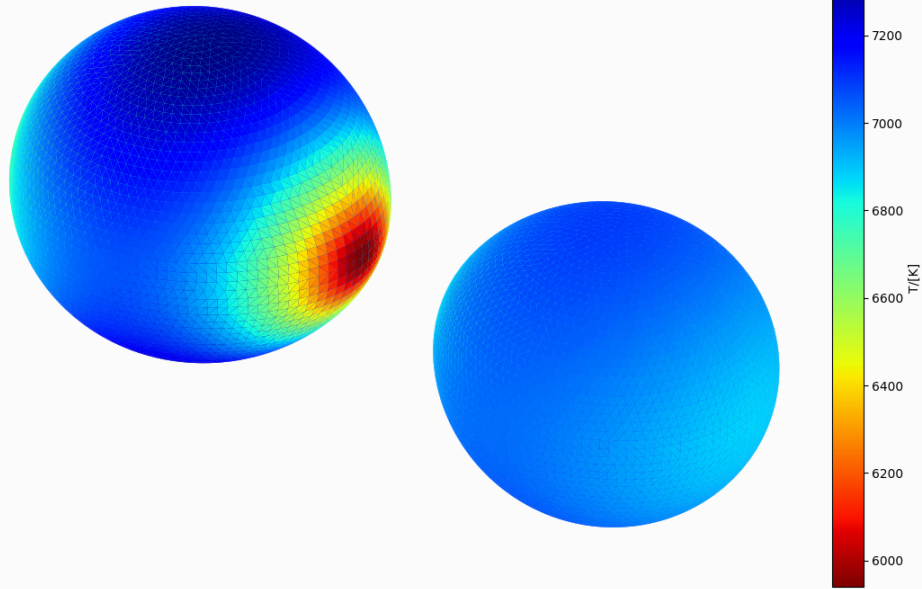
Roche geometry

- **Solution: equipotential surfaces.**
- **Solved iteratively.**
- **Tidally deformed surfaces of the components due to the mutual gravitational interaction.**

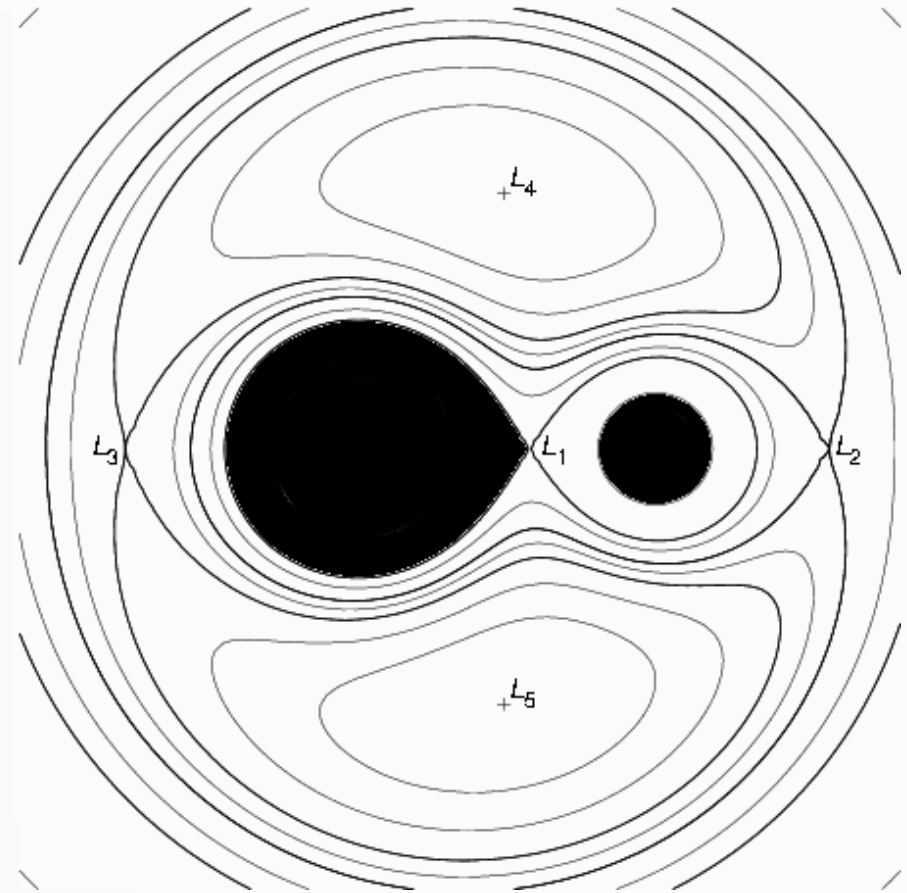
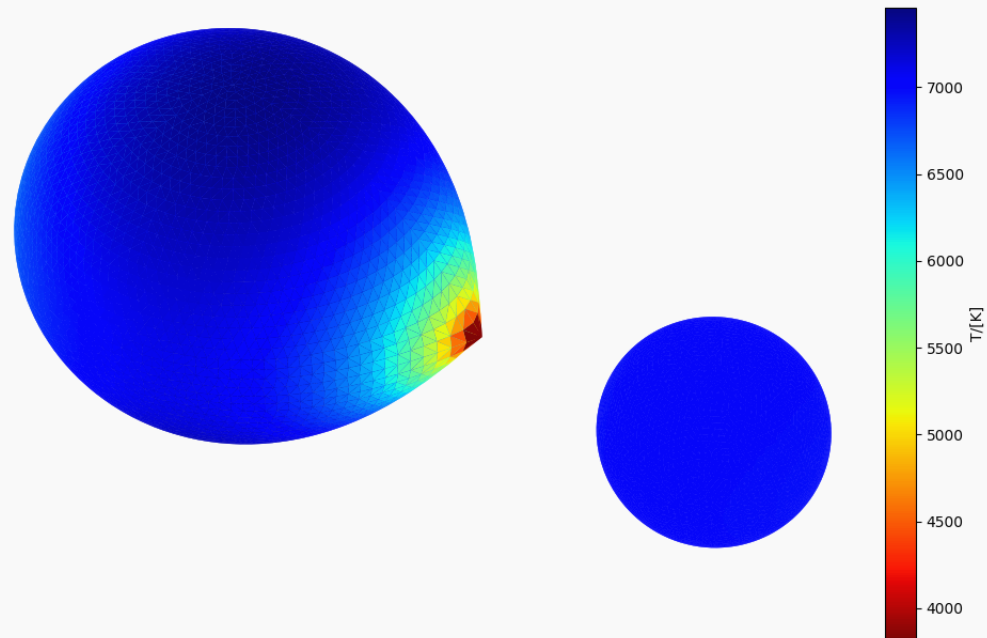


Source: Neslušan (2009)

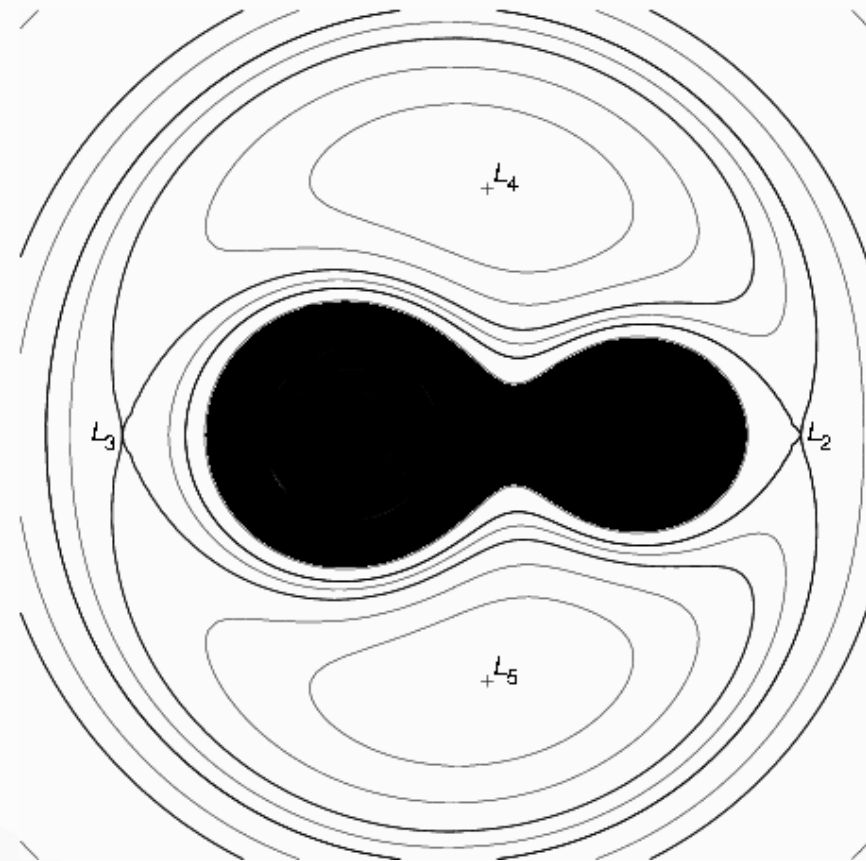
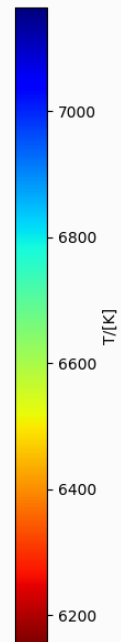
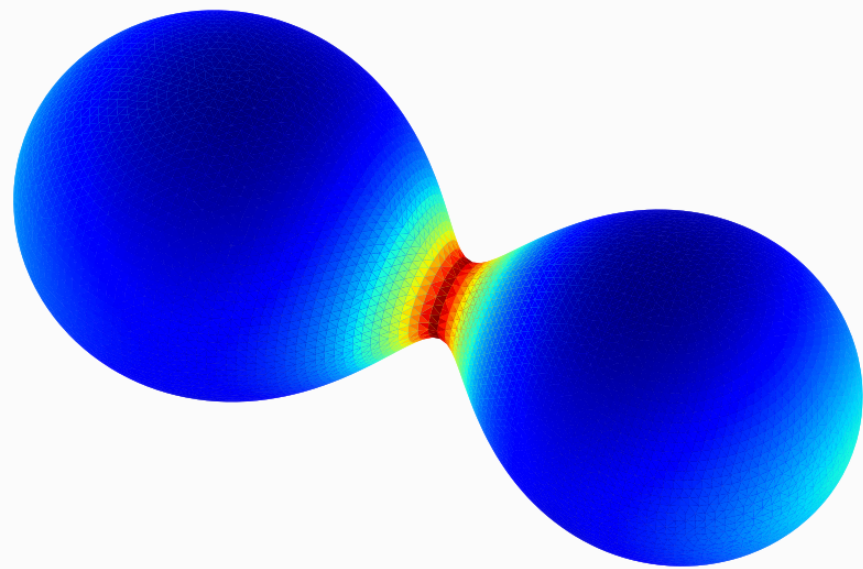
Detached binary



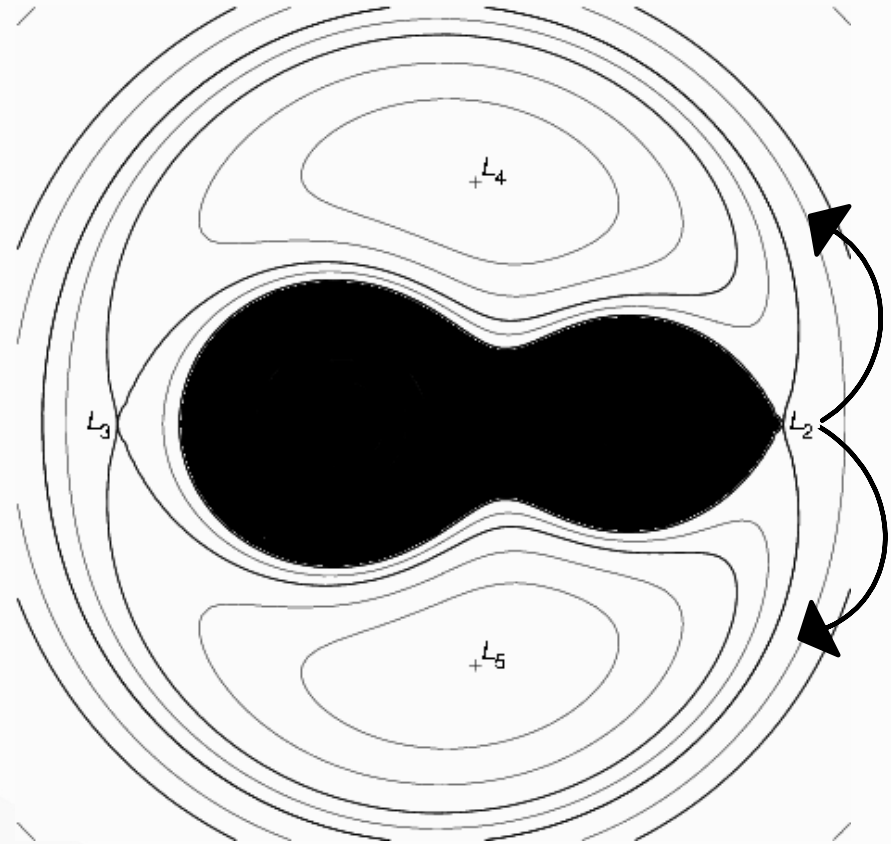
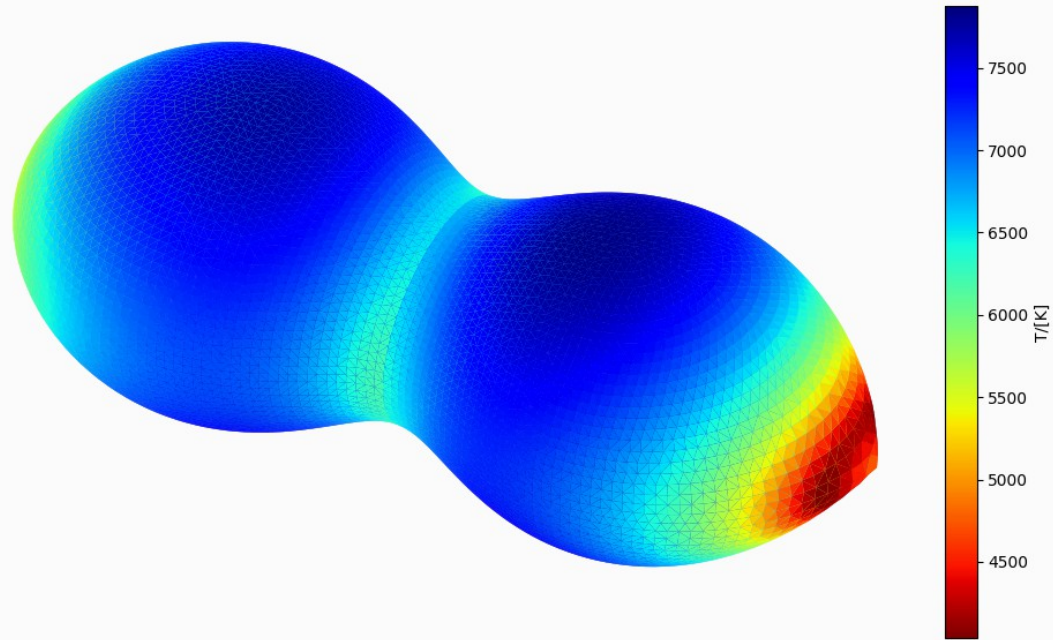
Semi-detached binary



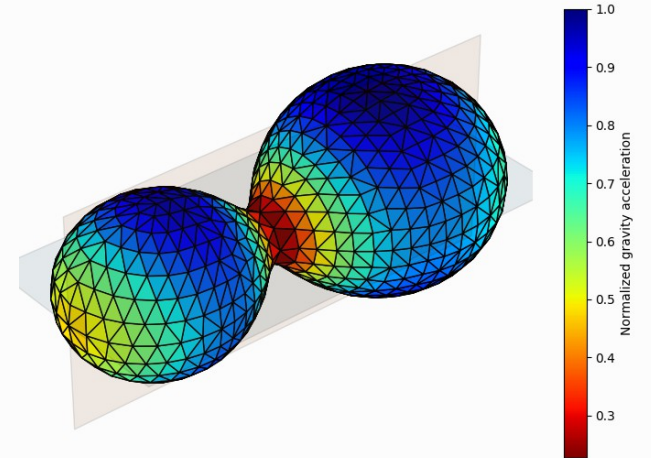
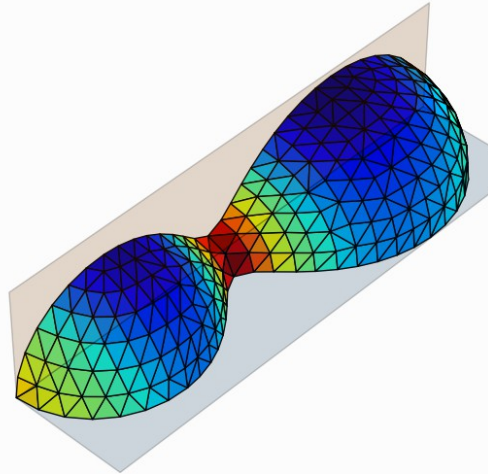
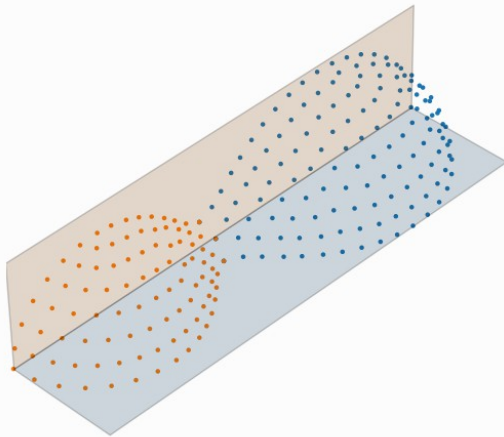
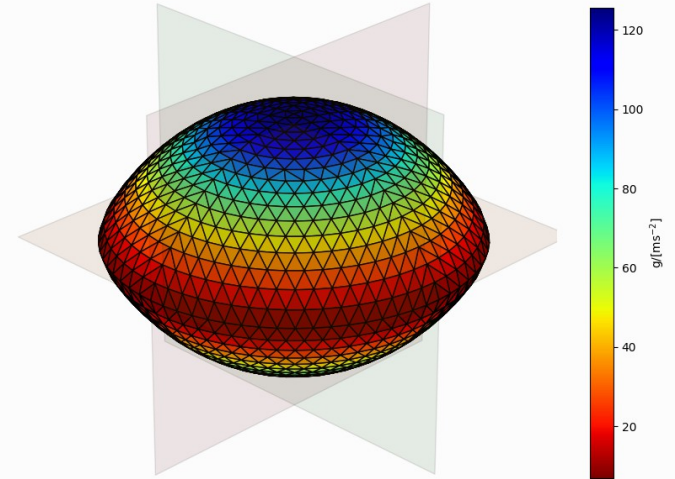
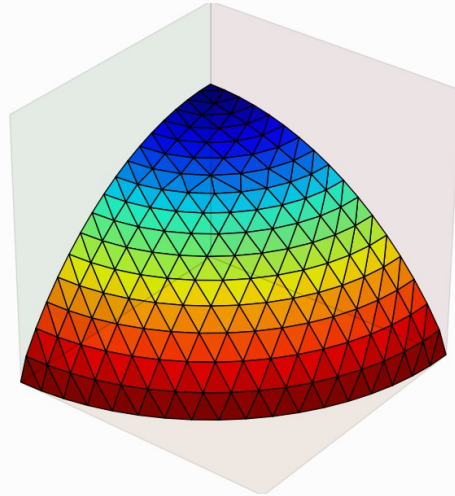
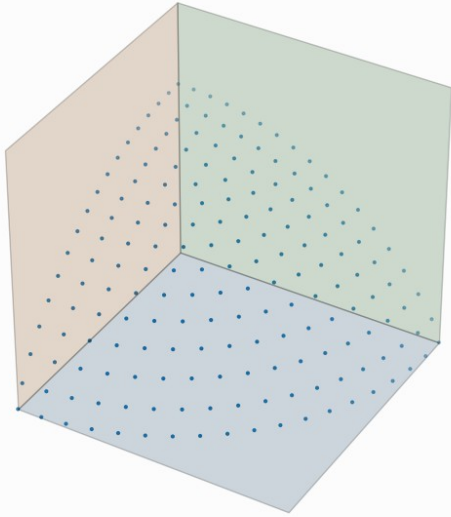
Over-contact binary



Over-contact binary with loss of material

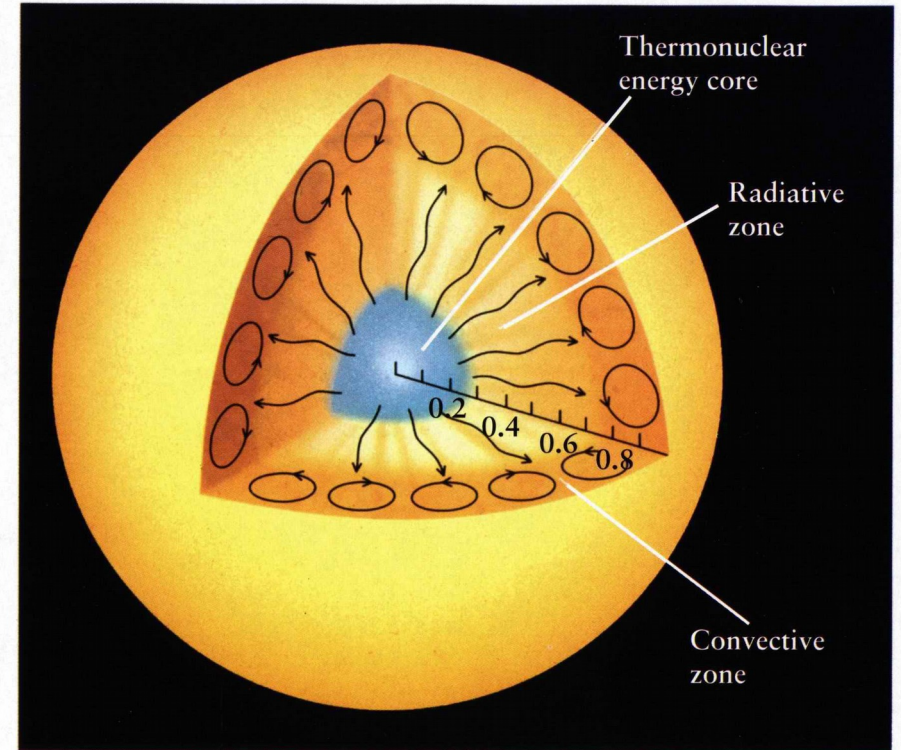


Symmetries of the stellar surfaces



Equilibrium stellar structure

- **Propagation of a stellar oscillations depends heavily on the internal stellar structure.**
- **Research on oscillations - window to internal stellar structure.**



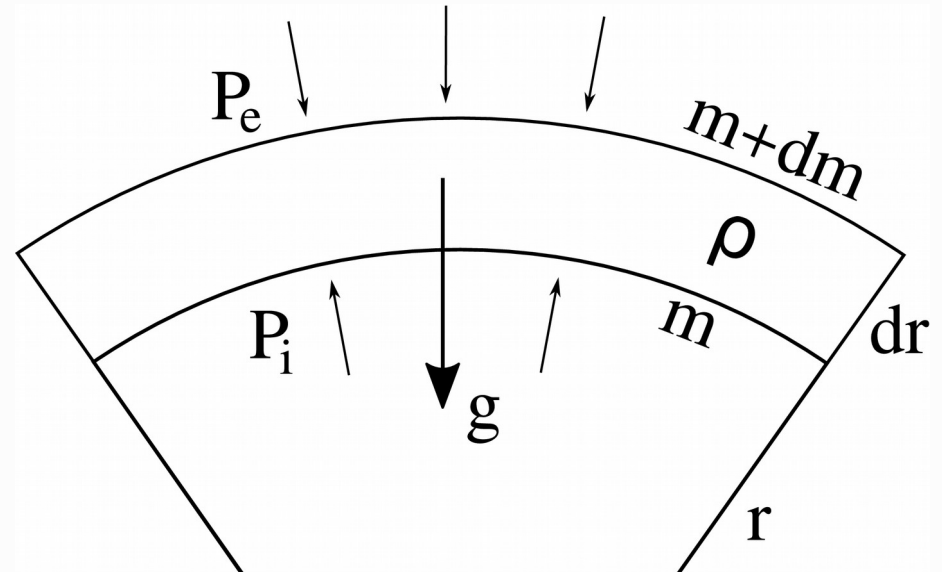
Equilibrium stellar structure

- **Maintained between gravity and internal pressure of material.**

$$\frac{\partial P}{\partial m} = -\frac{Gm}{4\pi r^4}$$

- **Switch from independent variable r to m .**

$$\frac{\partial r}{\partial m} = \frac{1}{4\pi r^2 \rho}$$



Equilibrium stellar structure – production of energy

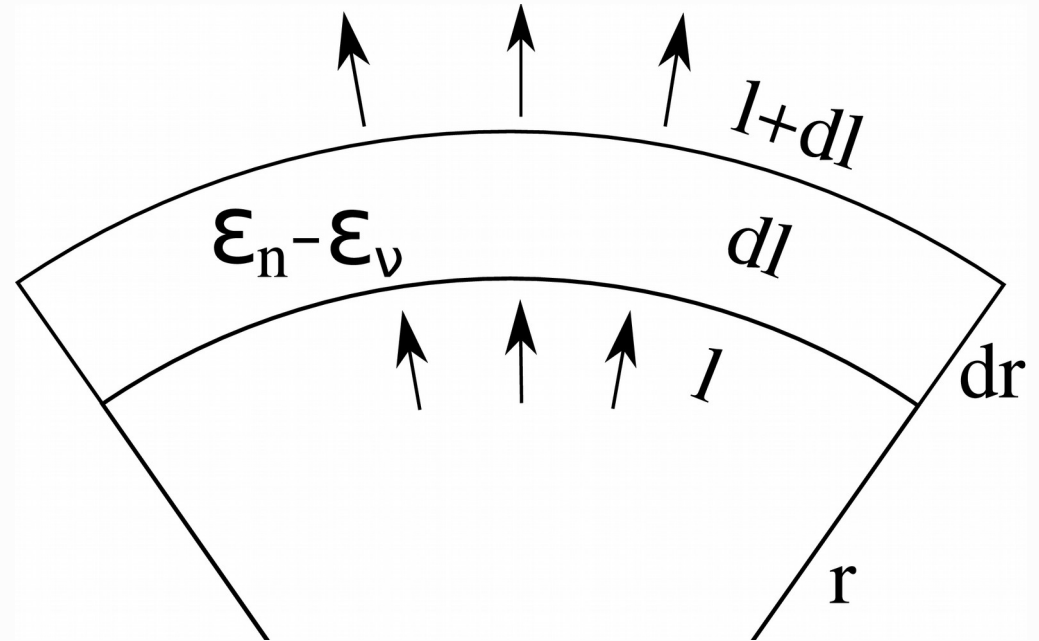
- **Distribution of luminosity l**

$$\frac{\partial l}{\partial m} = \epsilon_n - \epsilon_\nu$$

- **Energy produced via thermonuclear reactions.**

$$\epsilon = \sum_{ij} \epsilon_{ij} = \frac{1}{\rho} \sum_{ij} r_{ij} e_{ij}$$

$$\frac{\partial X_i}{\partial t} = \frac{m_i}{\rho} \left(\sum_j r_{ji} - \sum_k r_{ik} \right)$$



Equilibrium stellar structure – transport of energy

- **Radiation:**

- **temperature gradient necessary to transport all energy via radiation:**

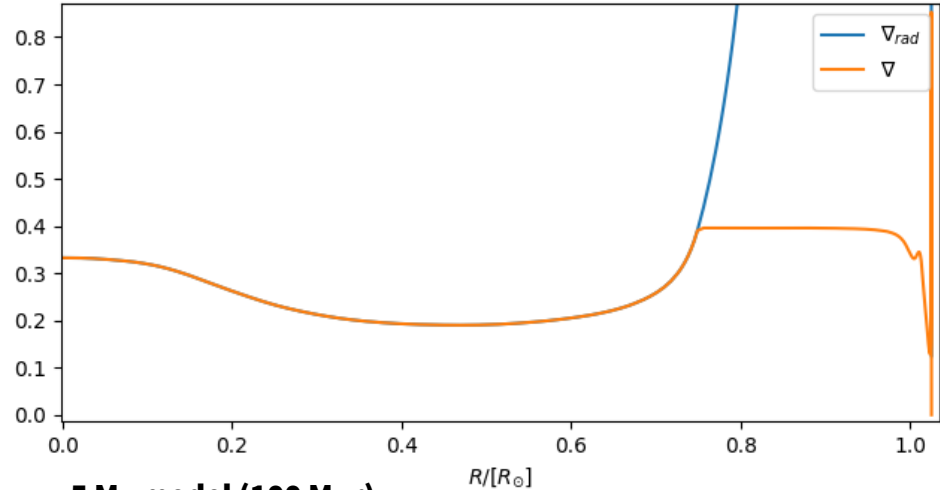
$$\nabla_{rad} = \left(\frac{\partial \ln T}{\partial \ln P} \right)_{rad}$$

- **Convection:**

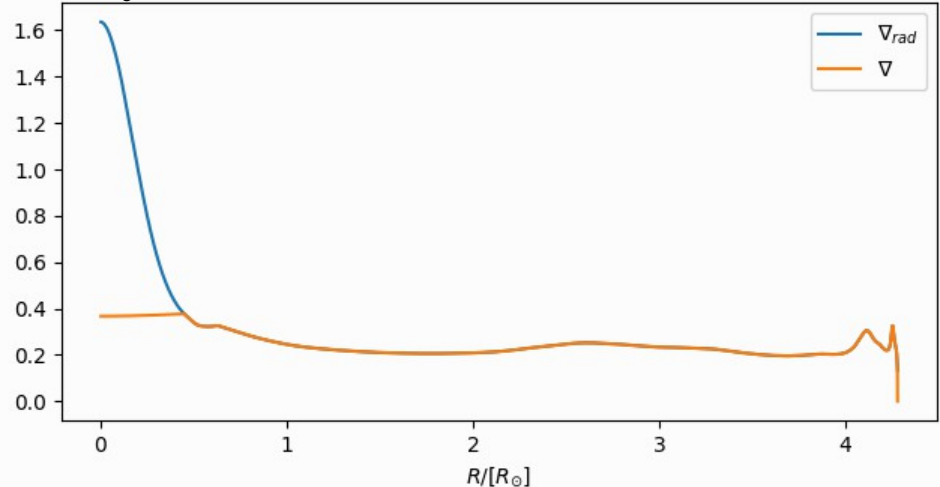
- **Fraction of the energy transported by convection if Ledoux criterion is violated:**

$$\nabla_{rad} < \nabla_{ad} + \frac{\delta}{\varphi} \frac{d \ln \mu}{d \ln P}$$

1 M_⊙ model (4.5 Gyr)



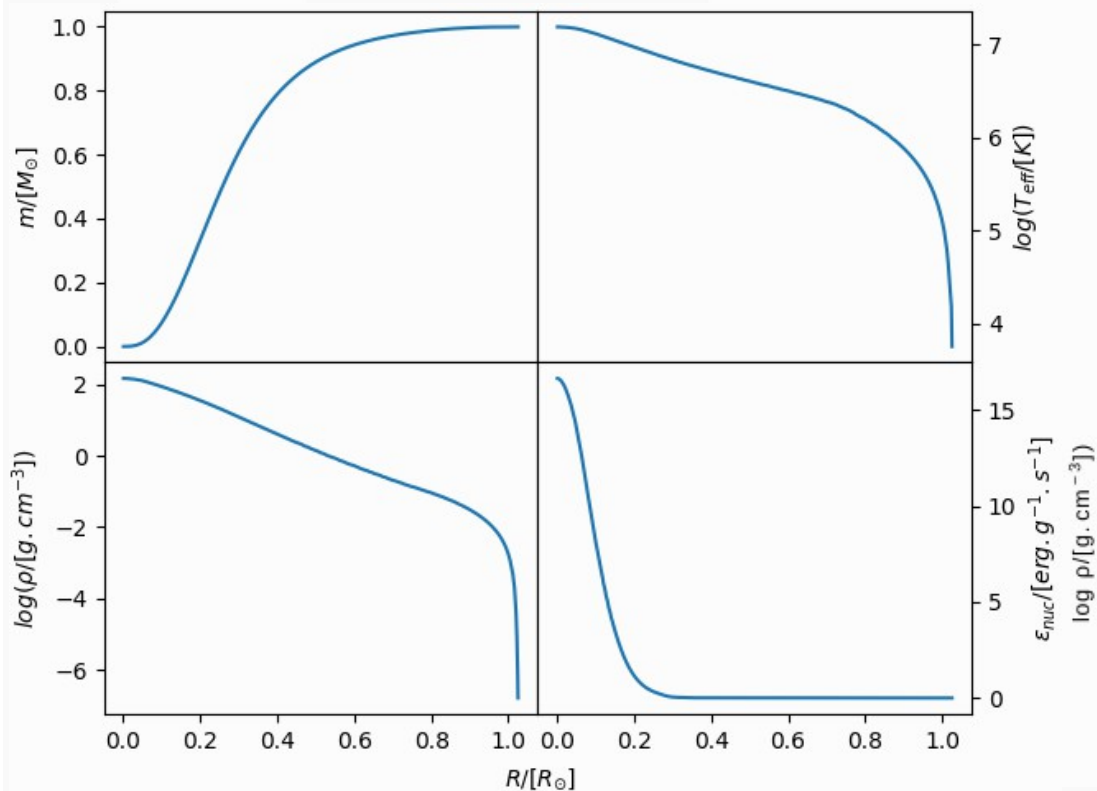
5 M_⊙ model (100 Myr)



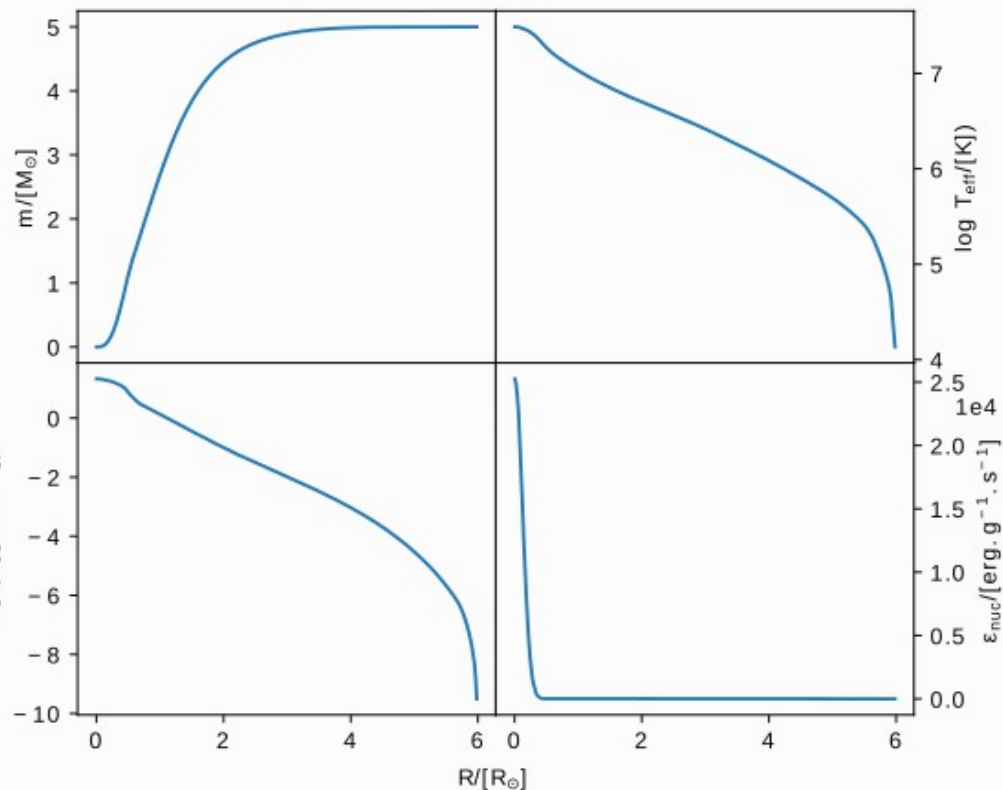
Equilibrium stellar structure

Modules for Experiments in Stellar Astrophysics (MESA)

- **1 M_{\odot} model (4.5 Gyr)**

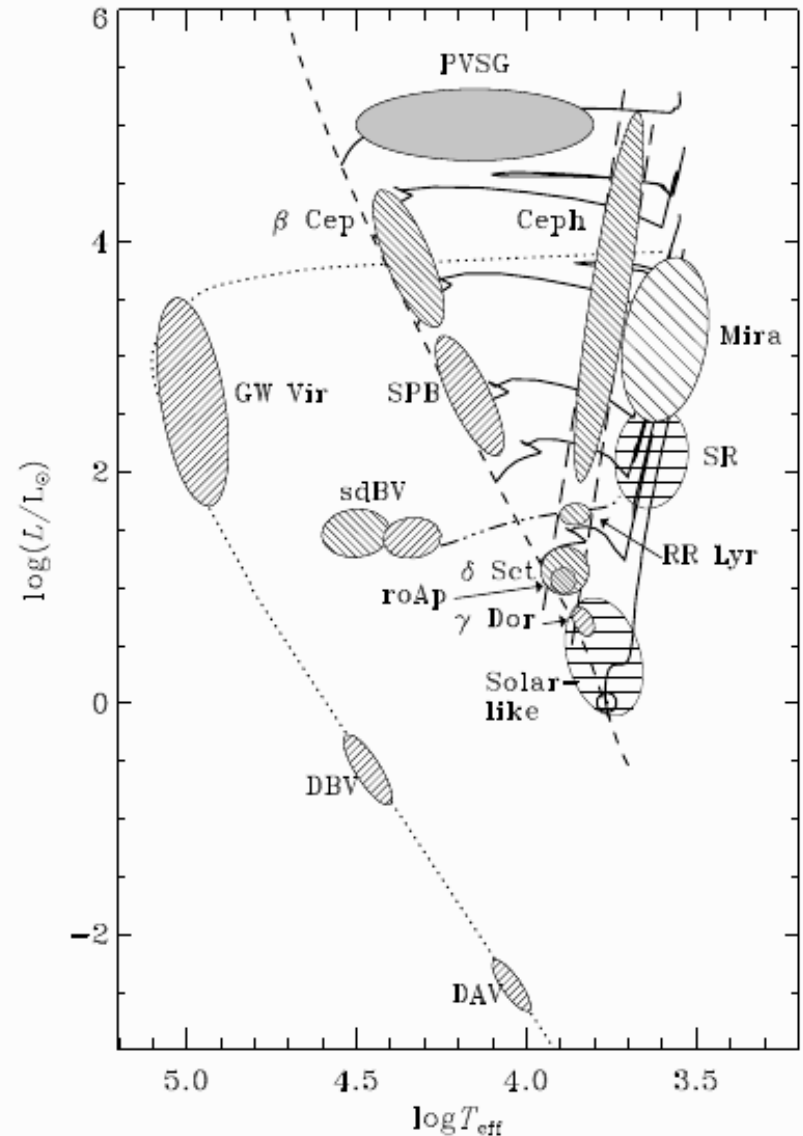


- **5 M_{\odot} model (100 Myr)**



Pulsations across the HR diagram

- **Periodic change in brightness.**
- **Change in luminosity due to internal processes.**
- **Multiple distinct groups distributed across the main sequence and the instability strip.**



Source: Aerts et al. (2010)

Stellar oscillations

- **Equilibrium disturbed by excitation mechanisms (κ -mechanism, convective blocking, tidal forces).**
- **Governed by hydrodynamic equations.**
- **Stellar matter modelled as non-viscous fluid.**

$$\frac{\partial \rho}{\partial t} + \nabla_r(\rho u) = 0$$

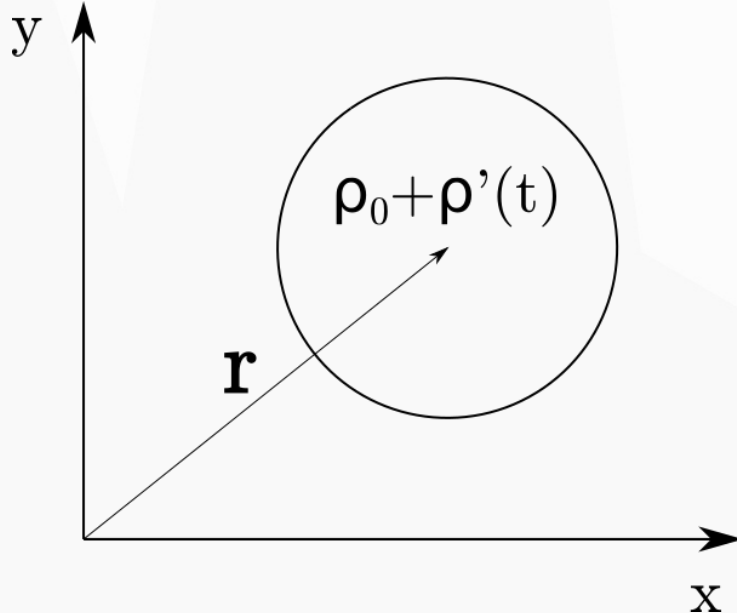
$$\left(\frac{\partial}{\partial t} + u \cdot \nabla_r \right) u = -\nabla_r \Psi - \frac{1}{\rho} \nabla_r P$$

$$\frac{dq}{dt} = \frac{1}{\rho(\Gamma_3 - 1)} \left(\frac{dP}{dt} - \frac{\Gamma_1 p}{\rho} \frac{d\rho}{dt} \right)$$

Stellar oscillations – perturbative approach

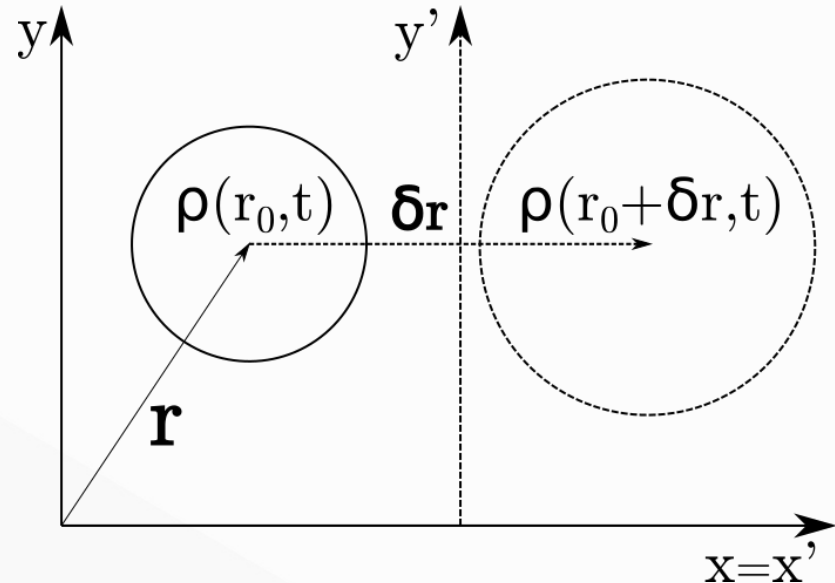
- **Oscillations treated as perturbation of equilibrium stellar structure.**
- **Eulerian perturbation**

$$\rho'(\mathbf{r}, t) = \rho(\mathbf{r}, t) - \rho_0(\mathbf{r})$$



- **Lagrangian perturbation**

$$\delta\rho(\mathbf{r}, t) = \rho'(\mathbf{r}_0, t) + \delta\mathbf{r} \cdot \nabla_{\mathbf{r}}\rho_0(\mathbf{r}_0).$$



Stellar oscillations – perturbative approach

- **Perturbative approach:**

- **Decomposition of stellar parameters to equilibrium and perturbed part.**
- **Subtraction of equilibrium part.**
- **Linearization.**

$$\rho' + \nabla_r \cdot (\rho_0 \delta \mathbf{r}) = 0$$

$$\rho_0 \frac{\partial^2 \delta \mathbf{r}}{\partial t^2} = -\rho_0 \nabla_r \Psi' - \rho' \nabla_r \Psi_0 - \nabla_r P'$$

$$P' + \nabla_r P_0 \cdot \delta \mathbf{r} = \frac{\Gamma_{1,0} P_0}{\rho_0} (\rho' + \nabla_r \rho_0 \cdot \delta \mathbf{r})$$

$$P' = \frac{k_B}{\mu} (\rho_0 T' + T_0 \rho')$$

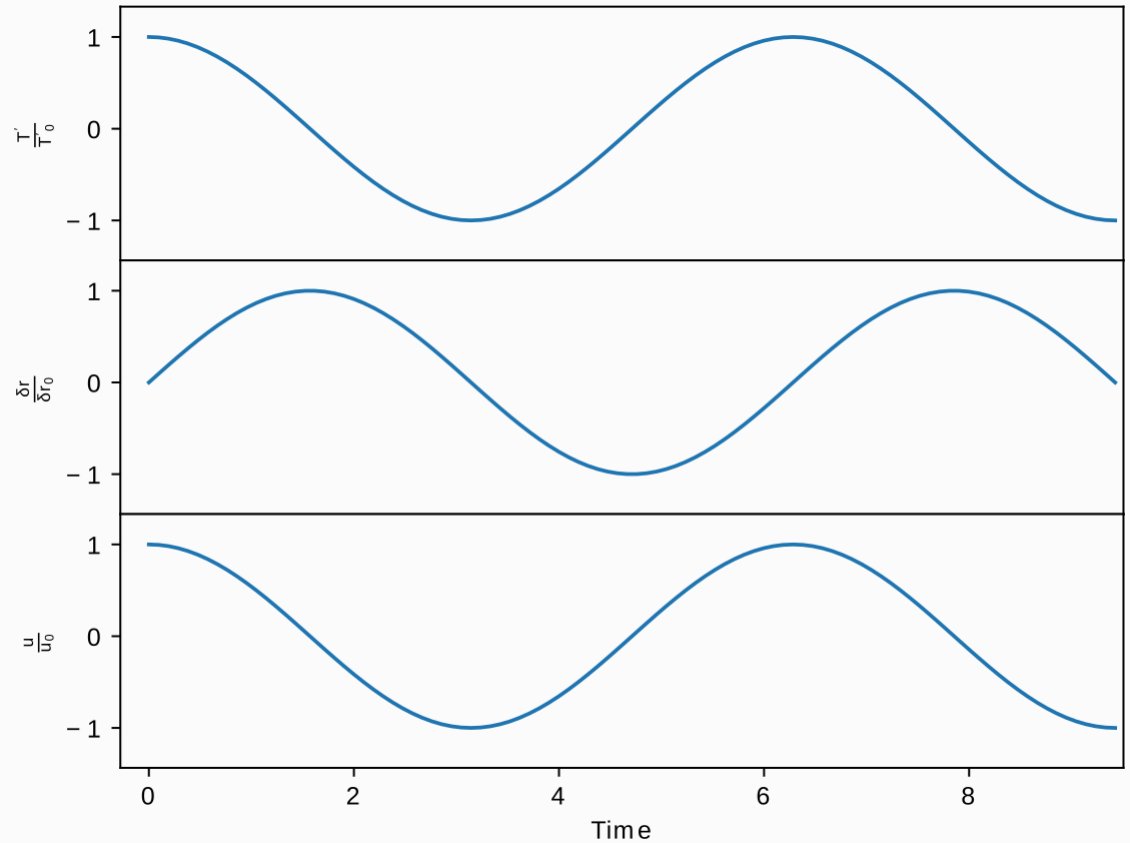
$$\Delta \Psi' = 4\pi G \rho'$$

Pressure driven oscillations

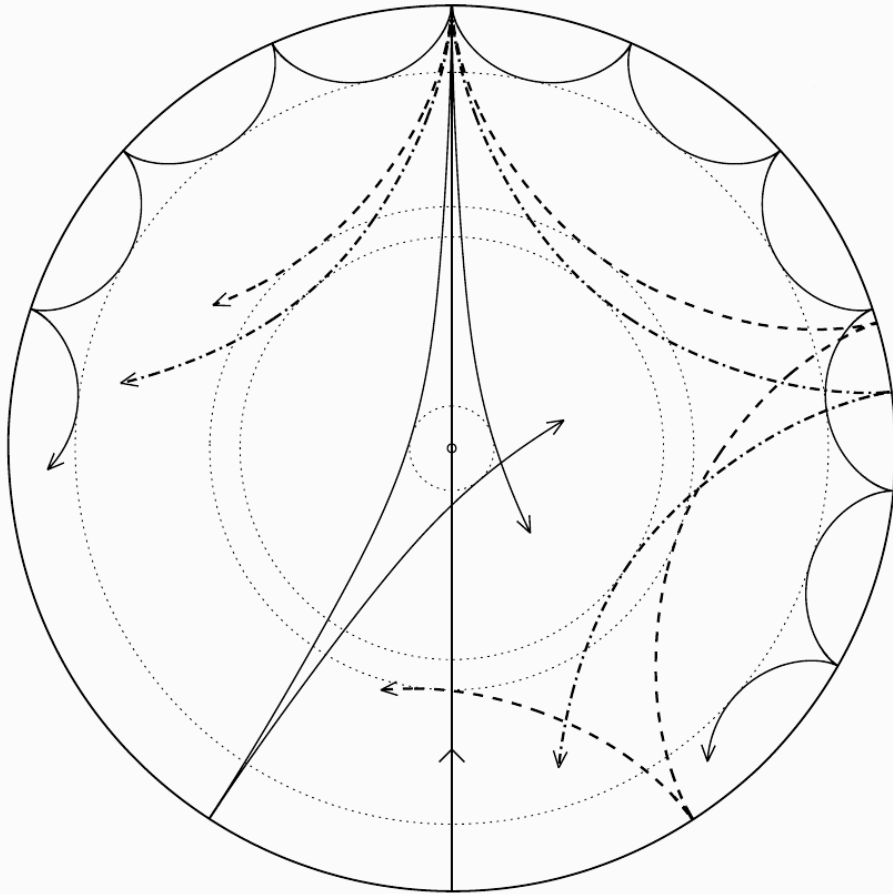
- **Acoustic waves**
- **Local approach – planar wave**

$$\frac{\partial^2 P'}{\partial t^2} = \frac{\Gamma_{1,0} P_0}{\rho_0} \Delta P'$$

$$c_0 = \sqrt{\frac{\Gamma_{1,0} P_0}{\rho_0}}$$

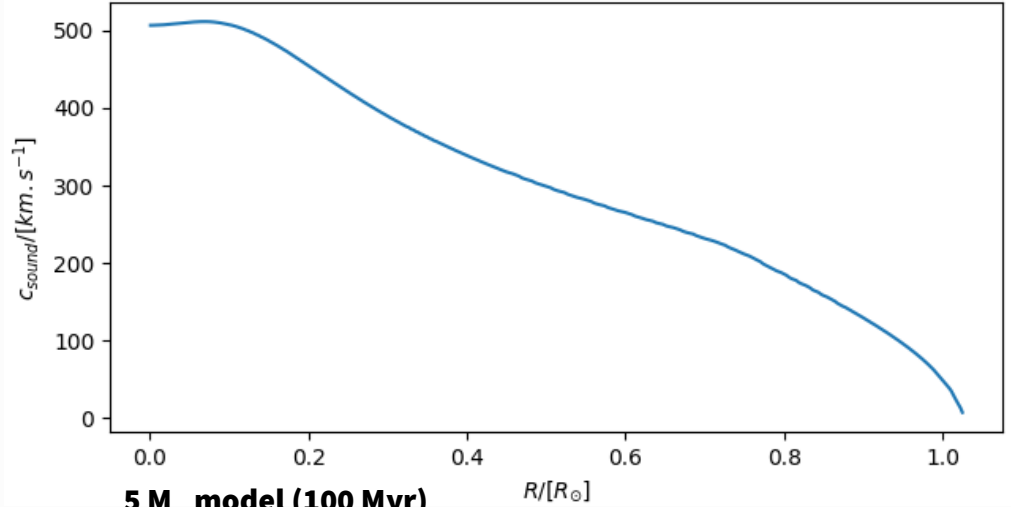


Pressure driven oscillations

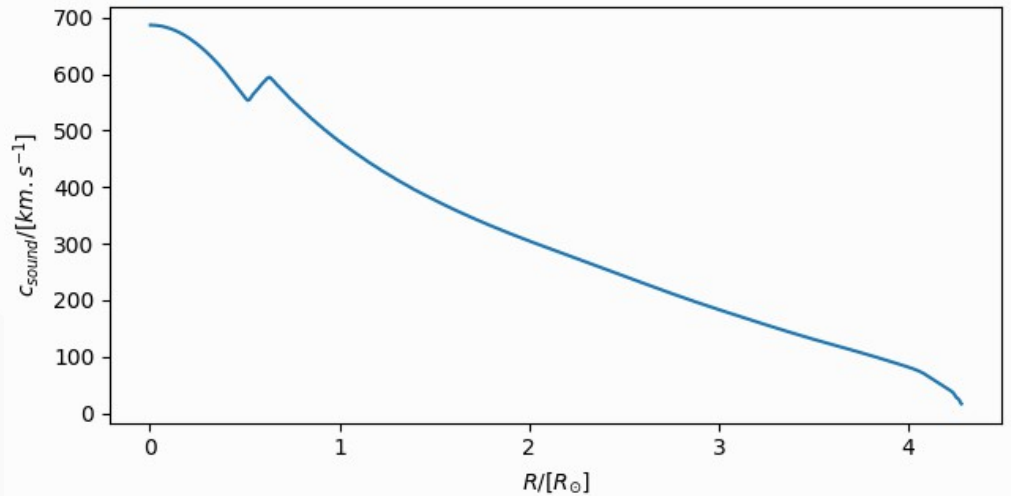


Source: Aerts et al. (2010)

1 M_{\odot} model (4.5 Gyr)

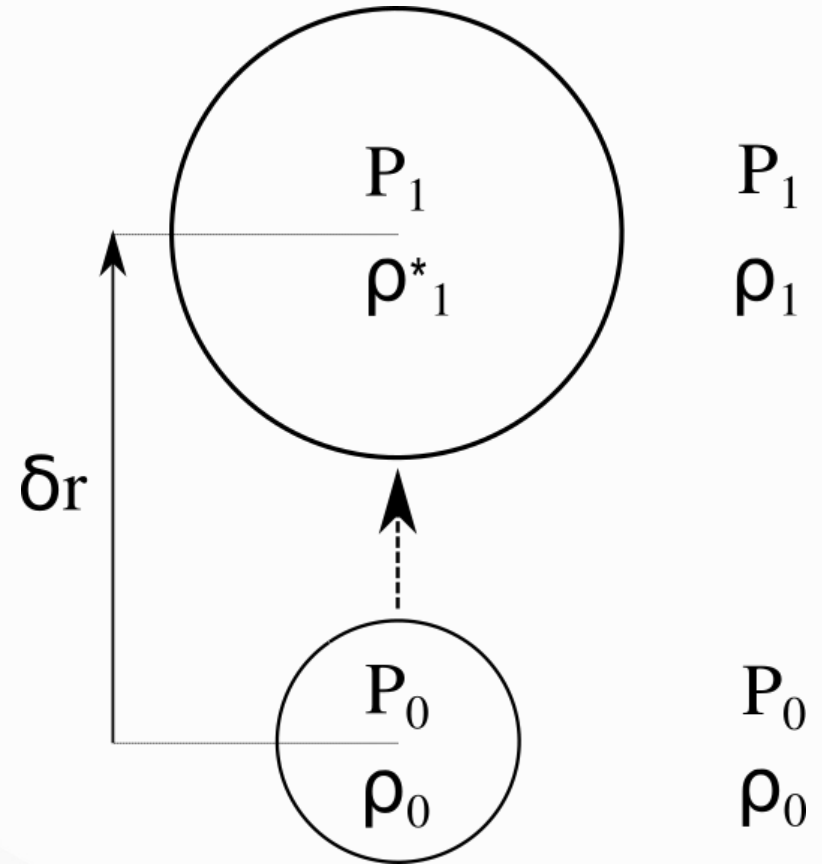
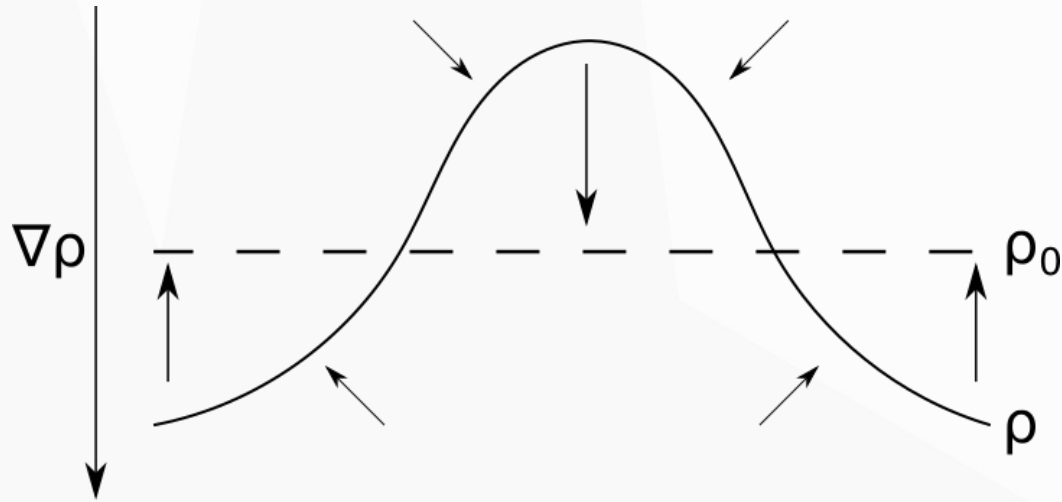


5 M_{\odot} model (100 Myr)



Gravity driven oscillations

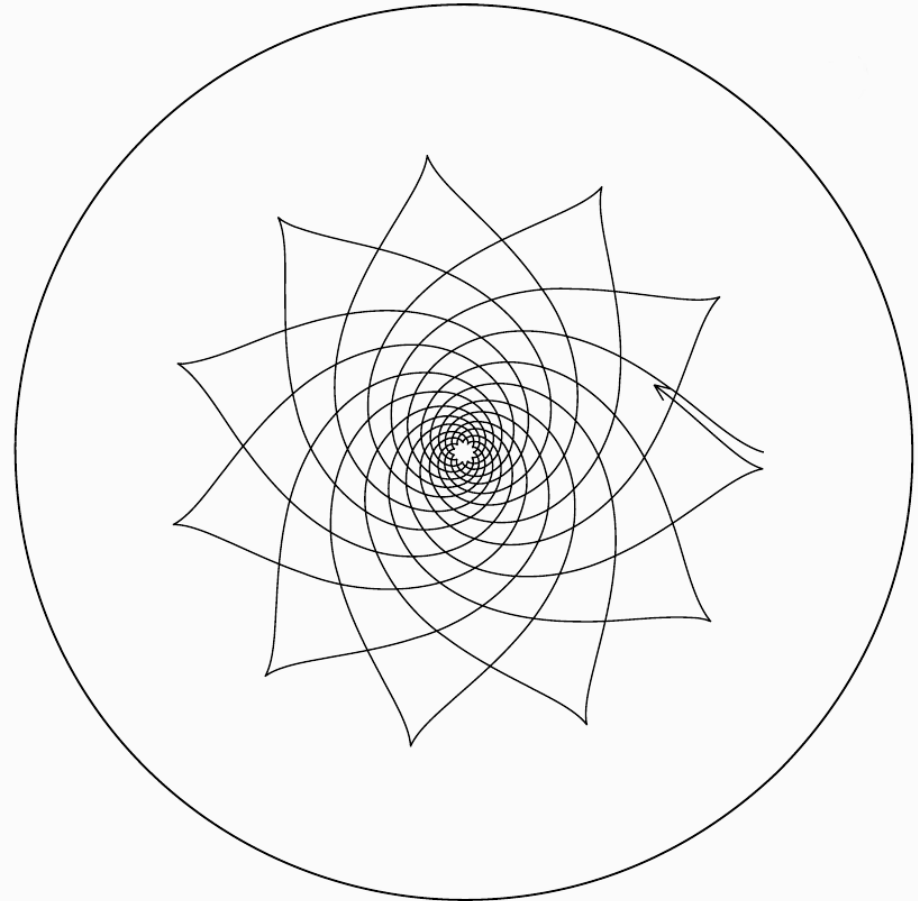
$$N^2 = \left(\frac{1}{\Gamma_1} \frac{d \ln P_0}{dr} - \frac{d \ln \rho_0}{dr} \right) g_0 > 0$$



Gravity driven oscillations

- **Dispersion relation**

$$\omega^2 = \frac{N^2}{1 + \left(\frac{k_r}{k_h}\right)^2}$$



Source: Aerts et al. (2010)

Surface representation of oscillations

- **Surface modes – spherical harmonics**

$$Y_l^m(\theta, \varphi) = (-1)^m c_{lm} P_l^m(\cos \theta) e^{im\varphi}$$

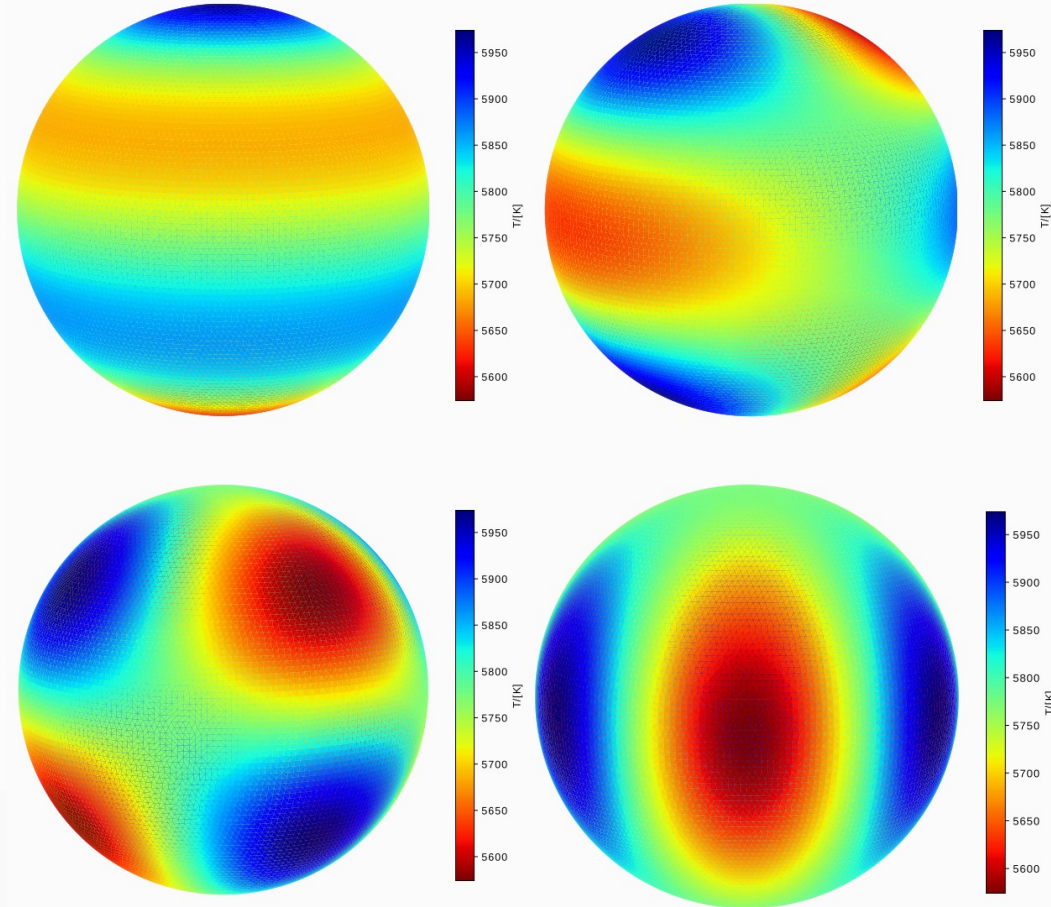
- **l – no. of surface nodal lines**

- **m – no. of azimuthal nodal lines, $|m| < l$**

$$\xi_r = \sqrt{4\pi} \tilde{\xi}_r(r) Y_l^m(\theta, \phi) \exp(-i\omega t),$$

$$T' = \sqrt{4\pi} \tilde{T}'(r) Y_l^m(\theta, \varphi) \exp(-i\omega t)$$

$$\tilde{T}'(r) = \frac{\tilde{P}'(r)\mu}{\rho_0 k_B} \left(\frac{\Gamma_{1,0} - 1}{\Gamma_{1,0}} \right)$$



Future goals of the project

- **Developing a code able to generate a light curves of eclipsing binaries with a surface inhomogenities.**
- **Explore the solutions of stellar oscillations for heavily distorted components.**
- **Acceleration of computations utilizing GPU computing.**
- **Adapting 1D stellar models to model interiors of deformed stellar components.**



Thank you!
(Any questions?)